# ABSTRACT OF THE TALK "DEFORMATION OF $G$-STRUCTURES" 

M. MALAKHALTSEV

## REFERENCES

Kodaira K., Complex manifolds and deformation of complex structures, 1981.
Pommaret, J.F. Systems of partial differential equations and Lie pseudogroups, Math. and Appl., 14 (1978).

## 1. G-STRUCTURES

Let $G$ be a subgroup in $G L(n)$. A $G$-structure on a manifold $M$ is a principal $G$-subbundle $L_{G}(M) \rightarrow M$ of the linear frame bundle $L(M)$. Each $G$-structure is defined by a section $s: M \rightarrow$ $P_{G}=L(M) / G$.

A $G$-structure on a manifold $M$ is mathcalled integrable if, for each point $p \in M$, there exists a chart on $U \ni p$ whose natural frame is a section of $L_{G}(M)$.

Example 1. Let

$$
G=\left(\begin{array}{cc}
A & B  \tag{1}\\
0 & C
\end{array}\right)
$$

where $A$ is $k \times k$-matrix, $B$ is $k \times(n-k)$-matrix, and $C$ is $(n-k) \times(n-k)$-matrix.
Then a $G$-structure $L_{G}(M) \rightarrow M$ is a $k$-dimensional distribution $\Delta$ on $M$. The corresponding section $s: M \rightarrow P_{G}=L(M) / G$ is a section of the Grassmann bundle $G_{k}(M)$ of $k$-dimensional subspaces in $T M$.

If this $G$-structure is integrable, then the distribution $\Delta$ is also integrable, hence determines a foliation.

## 2. Deformations of integrable $G$-Structures

2.1. Deformation of $G$-structures. A deformation of $a G$-structure $s$ is a one-parametric family of sections $s_{t}: M \rightarrow L(M) / G$ such that $s_{0}=s$. A deformation is mathcalled inessential if $s_{t}=f_{t}^{*}(s)$ for a one-parametric family $f_{t}$ of diffeomorphisms of $M$.

An infinitesimal deformation of a section $s: M \rightarrow P_{G}$ is a vertimathcal vector field $V=\left.\frac{d}{d t}\right|_{t=0} s_{t}$ along $s(M)$, or a section of the bundle $s^{*} V\left(P_{G}\right)$. That is

$$
\begin{equation*}
\mathcal{D}(s)=\Gamma\left(M ; s^{*} V\left(P_{G}\right)\right) \tag{2}
\end{equation*}
$$

An inessential infinitesimal deformation of a section $s: M \rightarrow P_{G}$ is a vertimathcal vector field $V=\left.\frac{d}{d t}\right|_{t=0} f_{t}^{*}\left(s_{0}\right)$ along $s(M)$. One can prove that the space of inessential infinitesimal deformations is

$$
\begin{equation*}
\mathcal{D}_{0}=\left\{L_{X} s \mid X \in \mathfrak{X}(M)\right\} \tag{3}
\end{equation*}
$$

Hence the space of essential deformations of $s$ is

$$
\begin{equation*}
\mathcal{D}_{e s s}=\frac{\mathcal{D}(s)}{\mathcal{D}_{0}(s)}=\frac{\Gamma\left(M ; s^{*} V\left(P_{G}\right)\right)}{\left\{L_{X} s \mid X \in \mathfrak{X}(M)\right\}} \tag{4}
\end{equation*}
$$

Remark 1. This can be done for sections of natural bundle.
2.2. Deformations of integrable structures. A deformation of an integrable $G$-structure $s$ is a one-parametric family of sections $s_{t}: M \rightarrow L(M) / G$ such that each $s_{t}$ is integrable and $s_{0}=s$.

An infinitesimal deformation of a section $s: M \rightarrow P_{G}$ is a vertimathcal vector field $V=\left.\frac{d}{d t}\right|_{t=0} s_{t}$ along $s(M)$, or a section of the bundle $s^{*} V\left(P_{G}\right)$ but now the sections $s_{t}$ are integrable $G$-structures.

First approach. Any integrable $G$-structure $s$ on $M$ determines a pseudogroup structure on $M$.
$\Gamma$ is a pseudogroup of transformations of $f: \mathbb{R}^{n} \rightarrow \mathbb{R}$ such that $D f$ lies in $G$. Then existence of $G$-structure is equivalent to existence of a $\Gamma$-atlas (an atlas on $M$ with transition functions in $\Gamma$ ).

We can repeat the construction of deformations of complex structure.
We introduce the sheaf $\mathfrak{X}_{a}$ of infinitesimal automorphisms of $s$. Then the vector fields

$$
\begin{equation*}
V_{\beta \alpha}=\left.\frac{d}{d t}\right|_{t=0} \varphi_{\beta}^{-1} \varphi_{\beta \alpha}(t) \varphi_{\alpha} \tag{5}
\end{equation*}
$$

give us a cocycle in $\check{C}^{1}\left(M ; \mathfrak{X}_{a}\right.$.
If the $\varphi_{\beta \alpha}(t)$ correspond to an unessential deformation, then $V_{\beta \alpha}=W_{\beta}-W_{\alpha}$, where $W_{\alpha} \in \mathfrak{X}_{a}\left(U_{\alpha}\right)$, so the cocycle is a coboundary. Thus

$$
\begin{equation*}
\mathcal{D}_{\text {ess }}=\check{H}^{1}\left(M ; \mathfrak{X}_{a}\right) \tag{6}
\end{equation*}
$$

Example 2. If $s$ is foliation, then the transition functions have the form:

$$
\begin{equation*}
\bar{x}^{i}=\varphi^{i}\left(x^{j}\right), \quad \bar{x}^{\alpha}=\varphi^{\alpha}\left(x^{j}, x^{\beta}\right) \tag{7}
\end{equation*}
$$

and the sheaf of infinitesimal isometries of a foliation $\mathcal{F}$ consists of so-called basic vector fields of $\mathcal{F}$ :

$$
\begin{equation*}
V^{i}\left(x^{j}\right) \partial_{i}+V^{\alpha}\left(x^{j}, x^{\beta}\right) \partial_{\alpha} \tag{8}
\end{equation*}
$$

Second approach A section $s: M \rightarrow P_{G}$ can be represented by a covering $U_{\alpha}$ such that on each $U_{\alpha}$ there is given a local frame $e_{\alpha}$ and $e_{\alpha}=e_{\beta} g_{\beta \alpha}$, where $g_{\beta \alpha} \in G$. If $s$ is an integrable $G$-structure, then the frames $e_{\alpha}$ are natural frames of charts on $M$.

The vector bundle $s^{*}\left(V P_{G}\right)$ is isomorphic to $F_{G}=T_{1}^{1}(M) / E_{G}$, where $E_{G}$ are linear operators whose matrices with respect to $e_{\alpha}$ lie in $\mathfrak{g}$.

It will be convenient to see it in the following way. Let $\partial_{i}$ be the frame determining the section $s$, and $e_{i}(t)=A_{i}^{j}(t) \partial_{j}$ be the frame which determines the deformation. Then $\left.\frac{d}{d t}\right|_{t=0} e_{i}=\left.\frac{d}{d t}\right|_{t=0} A_{i}^{j} \partial_{j}$, so we have linear operators $V_{i}^{j}=\left.\frac{d}{d t}\right|_{t=0} A_{i}^{j}$. However, $\bar{e}_{i}(t)=g_{i}^{m}(t) A_{m}^{j}(t) \partial_{j}$ determine exactly the same deformation, and if we take the derivative we get that $\bar{V}_{i}^{j}=V_{i}^{j}+W_{i}^{j}$, where $W_{i}^{j}$ is in $E_{G}$.

Now, if we take deformation, consisting of integrable structures, then $\left[e_{i}(t), e_{j}(t)\right]=0$, from where we get that $\partial_{[i} V_{j]}^{k}=0$ (of course, this equality is not invariant). However, if we take into account that we are working with classes and use the torsion free connection $\nabla$ such that $\nabla s=0$, or equivalently, the connection form takes values in $\mathfrak{g}$, then we arrive at the result that the deformation space is

$$
\begin{equation*}
\mathcal{D}=\operatorname{ker}\left(D: \frac{\Omega^{1} M \otimes T M}{\Omega^{0} M \otimes E_{G}} \rightarrow \frac{\Omega^{2} M \otimes T M}{\operatorname{Alt}\left(\Omega^{1} M \otimes E_{G}\right)}\right. \tag{9}
\end{equation*}
$$

Example 3. For a foliation $\mathcal{F}$ we get that $\mathcal{D}=\left\{V_{\alpha}^{i}: T \mathcal{F} \rightarrow T M / T \mathcal{F} \mid \partial_{[\beta} V_{\alpha]}^{k}=0\right\}$.
Now if $s_{t}=f_{t}^{*}(s)$, where $f_{t}$ is a flow of a vector field $X$, is unessential deformation, then the vertical vector field $\left.\frac{d}{d t}\right|_{t=0}$ can be written in terms of $T_{1}^{1} / E_{G}$ as $\left[\partial_{i} X^{j}\right]$, so

$$
\begin{equation*}
\mathcal{D}_{\text {ess }}=\operatorname{ker}\left(D: \frac{\Omega^{1} M \otimes T M}{\Omega^{0} M \otimes E_{G}} \rightarrow \frac{\Omega^{2} M \otimes T M}{\operatorname{Alt}\left(\Omega^{1} M \otimes E_{G}\right)}\right) / \operatorname{im}\left(D: T M \rightarrow \frac{\Omega^{1} M \otimes T M}{\Omega^{0} M \otimes E_{G}}\right) \tag{10}
\end{equation*}
$$

Example 4. For a foliation $\mathcal{F}$ we get that $\mathcal{D}_{\text {ess }}=\left\{V_{\alpha}^{i}: T \mathcal{F} \rightarrow T M / T \mathcal{F} \mid \partial_{[\beta} V_{\alpha]}^{k}=0\right\} /\left\{\partial_{\alpha} X^{i}\right\}$.

## Relation between these approaches.

## 3. $P$-complex for the Lie derivative

For an involutive differential operator Spencer constructed a differential complex (Spencer, Pommaret).

Theorem 1. The P-complex of the Lie derivative is isomorphic to the complex $\left(C^{q}(P), d\right)$, where

$$
C^{q}(P)=\frac{\Omega^{q}(M) \otimes T M}{\operatorname{Alt}\left(\Omega^{q-1}(M) \otimes E_{\mathfrak{g}}\right)}
$$

and the differential $d: C^{q}(P) \rightarrow C^{q+1}(P)$ is induced by the differential operator $D=$ Alt $\circ \nabla$, where Alt is the alternation and $\nabla$ is the covariant derivative of a torsion-free connection adapted to $Q$ (i. e. $(D \omega)_{i_{1} \ldots i_{q+1}}^{j}=\nabla_{\left[i_{1}\right.} \omega_{\left.i_{2} \ldots i_{q+1}\right]}^{j}$ with respect to local coordinates adapted to $Q$ ).

Because, the kernel of the Lie derivative is exactly the sheaf $\mathfrak{X}_{a}$ of infinitesimal automorphisms of the $G$-structure $s$, we get a resolution for the $\mathfrak{X}_{a}$.

If this resolution is fine (the Poincaré Lemma holds true), then $\check{H}^{1}\left(M ; \mathfrak{X}_{a}\right) \cong H^{1}\left(C^{*}, d\right)$.
I do not know the general proof of the fact that this resolution is fine for arbitrary integrable $G$-structure, and possibly this is not true. However, for all the "classical" structures this resolution gives the corresponding "classical cohomology theory", and is fine.

Example 5. Let $\mathcal{F}$ be a foliation structure on a smooth manifold $M$, and let $\Delta$ be the corresponding integrable distribution. Then $E_{\mathfrak{g}}=\left\{A \in T_{1}^{1}(M) \mid A(\Delta) \subset \Delta\right\}$. Therefore $C^{p}=\Omega^{p}(\Delta) \otimes(T M / \Delta)$. If $\left(x^{i}, x^{\alpha}\right)$ are adapted local coordinates, i.e., if $\Delta$ is given by the equations $d x^{i}=0$, then $d$ can be written locally as $(d \omega)_{\alpha_{1} \ldots \alpha_{q+1}}^{i}=\partial_{\left[\alpha_{1}\right.} \omega_{\left.\alpha_{2} \ldots \alpha_{q+1}\right]}$. Thus we arrive at Vaisman's foliated cohomology.

